

# Charm quark physics from lattice QCD

arXiv:1104.4600,1111.0142

Yusuke Namekawa(Univ. of Tsukuba)  
for the PACS-CS collaboration

S.Aoki, K-I.Ishikawa, N.Ishizuka, T.Izubuchi, K.Kanaya, Y.Kuramashi, Y.Nakamura, Y.Namekawa,  
M.Okawa, Y.Taniguchi, A.Ukawa, N.Ukita, T.Yamazaki, T.Yoshie

# Contents

<b>1</b>	<b><u>Introduction</u></b>	<b>3</b>
<b>2</b>	<b><u>Simulation setup</u></b>	<b>7</b>
<b>3</b>	<b><u>Results</u></b>	<b>9</b>
3.1	<u>Charmonium spectrum</u> . . . . .	9
3.2	<u>Charm-strange spectrum</u> . . . . .	10
3.3	<u>Charm-ud spectrum</u> . . . . .	12
3.4	<u>Charm quark mass</u> . . . . .	13
3.5	<u>Decay constants and CKM matrix elements</u> . . . . .	14
3.6	<u>Charmed baryon</u> . . . . .	16
<b>4</b>	<b><u>Summary</u></b>	<b>18</b>

# 1 Introduction

[Charm quark system]

Charm quark system is charming for physicists.

- Charm quark mass and Cabibbo-Kobayashi-Maskawa(CKM) matrix elements are fundamental parameters of the standard model, which are hard to be determined by experiments.  
← In addition, these parameters are also needed as inputs for a new theory beyond the standard model.
- Exotic hadrons such as  $Z^+(4430)$ , made of  $udc\bar{c}$  !?, have been observed.
- Charmonium is a good probe for a hot and dense QCD matter.
  - ◇ Charm quark system is hard to be studied analytically because of  $m_{charm} \sim \Lambda_{QCD}$ . Effective theories are not effective. Non-perturbative method is needed.

## [Model and lattice QCD]

So far, many models have been used for study of charm quark system.

- Correctness of a model must be always checked, because the result is model-dependent.

← Experiment gives a check. In addition, precise lattice QCD calculations can judge a model now.

	Model	Lattice
Result	Model-dependent	Model-independent
Input	Many parameters	$\alpha_s, m_{quark}$ (or hadron masses)
$\alpha_s, m_{quark}$	Artificial	QCD running
Heavy quark	$1/M$ expansion	Full order
Cost	Low	High

## [Recent development of lattice QCD]

- Thanks to recent development of computers and algorithms, realistic lattice QCD simulations can be performed.  
→ Pion mass in lattice simulations reaches the physical value.

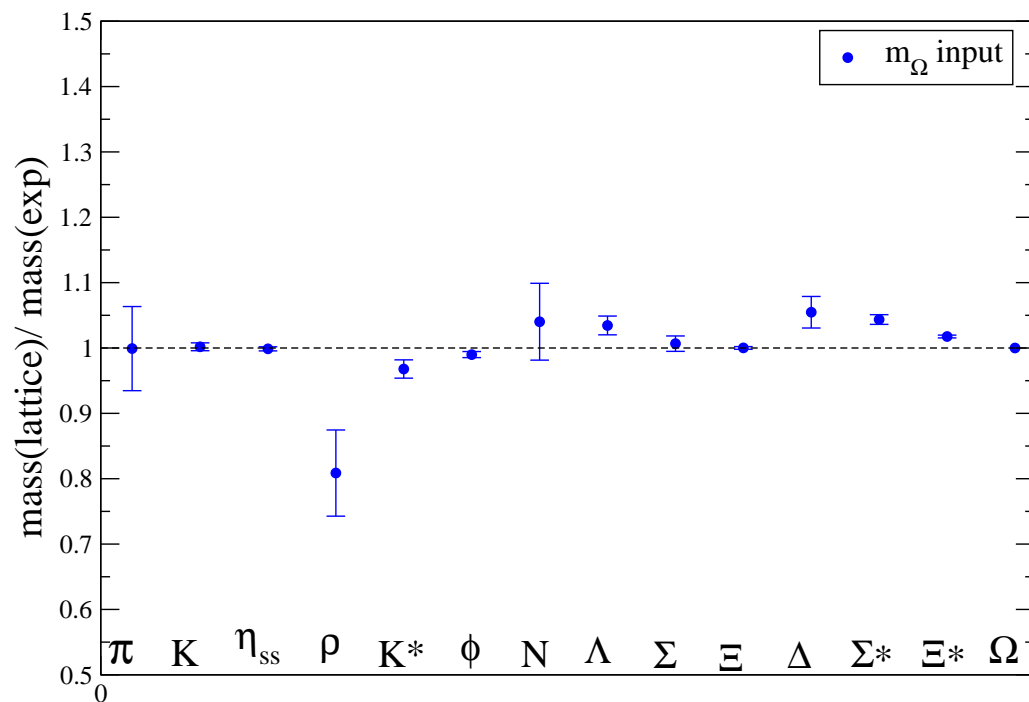
Year	Machine	Speed [TFlops]	$m_\pi$ [MeV]
1996-2005	CP-PACS	0.6	700
2006-	PACS-CS	14	160
2008-	T2K (Tokyo, Tsukuba, Kyoto)	235	135
Experiment			135



## [Light hadron spectrum]

Two lattice groups have reached the physical point of dynamical  $ud, s$  quarks. Light hadron spectrum is reproduced in 5% accuracy. PACS-CS,2009; BMW,2010  
→ Based on this result, we move on to the heavy quark system.

- For unstable hadrons such as  $\rho$ , more detailed analysis using Lüscher's formula is needed.
- MILC has started  $N_f = 2 + 1 + 1$  lattice QCD simulations around the physical point.



PACS-CS,2009

## 2 Simulation setup

We perform  $N_f = 2 + 1$  full QCD simulation (including dynamical up, down and strange quarks) for the charm quark system on the physical point.

- Action : Iwasaki gauge +  $O(a)$  improved Wilson fermion for light sea quarks + relativistic heavy fermion for valence charm quark
- Lattice size :  $32^3 \times 64$  ( $L = 3$  fm,  $a^{-1} = 2.2$  GeV ( $\beta = 1.90$ ))
- Sea and valence quark masses : on the physical point (i.e.  $m_\pi = 135$  MeV)
- Inputs :  $m_\pi, m_K, m_\Omega$  for  $m_{ud}, m_s, a$ ;  $m(1S) \equiv \frac{1}{4}(m_{\eta_c} + 3m_{J/\psi})$  for  $m_{charm}$

$m_{ud}^{\overline{\text{MS}}}(\mu = 2\text{GeV})[\text{MeV}]$	$m_s^{\overline{\text{MS}}}(\mu = 2\text{GeV})[\text{MeV}]$	$N_{conf}$ (MD time)
3	93	80 (2000)

## [Operators]

- We use the relativistic operators, because the relativistic heavy quark formulation is employed.
- We employ the two quark operators for mesons.  
← Only the quantum number is meaningful for the lattice field theory. Two- and four-quark operators give the same central value.

## [Operators for mesons]

$$\begin{aligned}M_{\Gamma}^{fg}(x) &= \bar{q}_f(x)\Gamma q_g(x), \\ \Gamma &= I, \gamma_5, \gamma_\mu, i\gamma_\mu\gamma_5, i[\gamma_\mu, \gamma_\nu]/2, \\ f, g &: \text{labels for quark flavors.}\end{aligned}$$

## [Operators for baryons with $J = 1/2$ ]

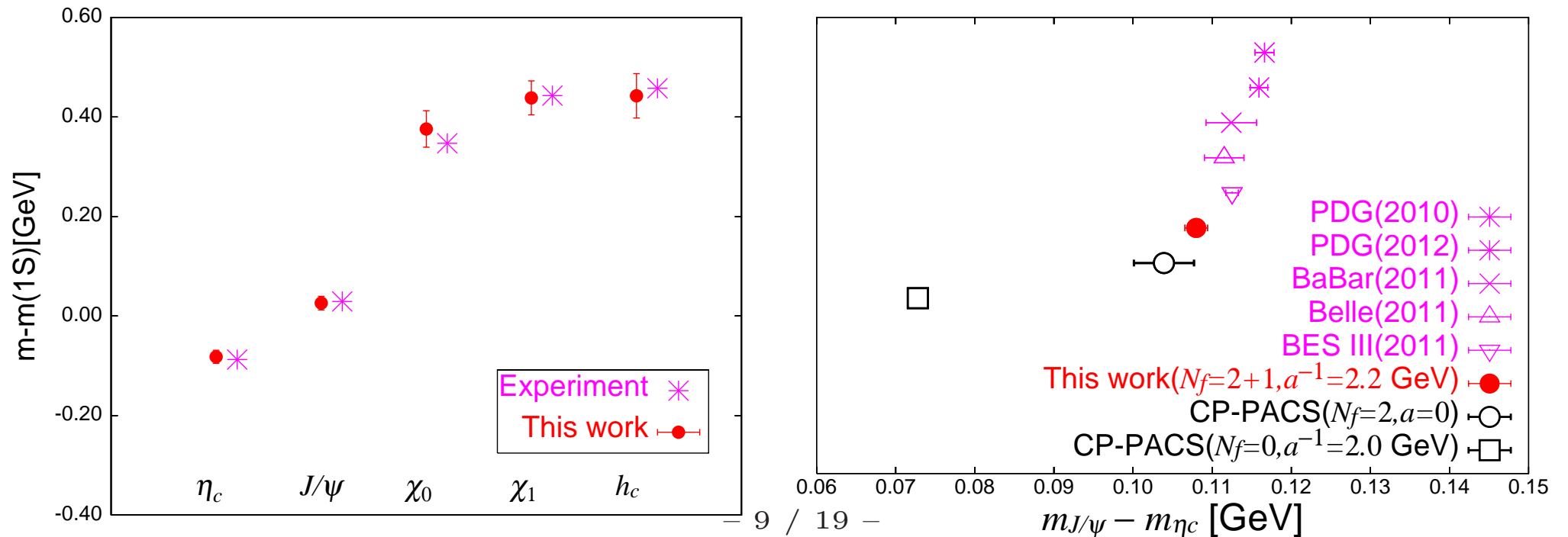
$$\begin{aligned}B_{\alpha}^{fgh}(x) &= \epsilon^{abc}((q_f^a(x))^T C\gamma_5 q_g^b(x))q_{h\alpha}^c(x), \\ C &= \gamma_4\gamma_2, \alpha = 1, 2.\end{aligned}$$



# 3 Results

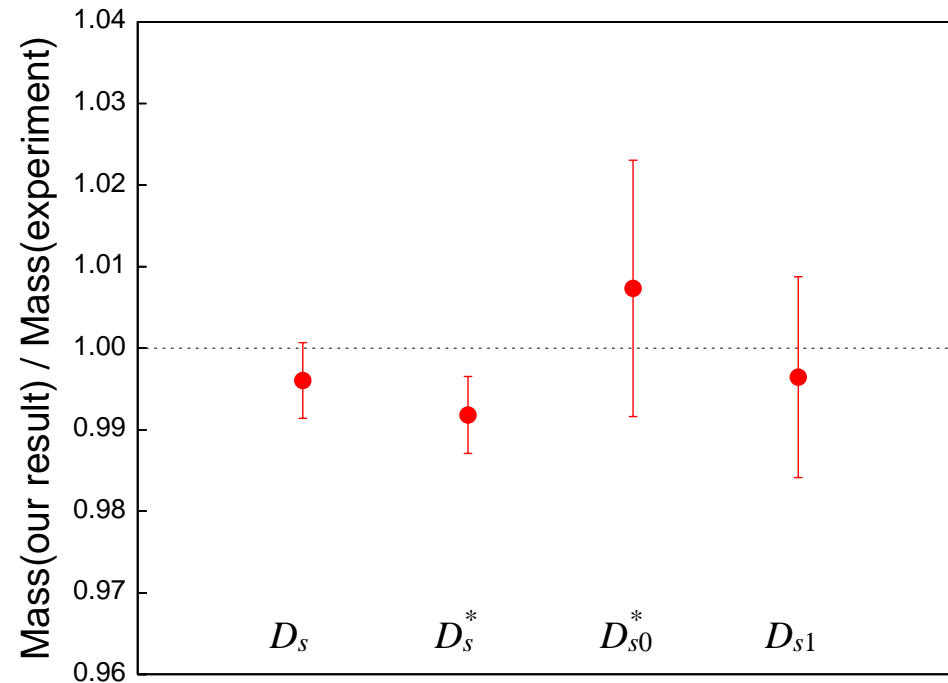
## 3.1 Charmonium spectrum

- Since  $m(1S) \equiv \frac{1}{4}(m_{\eta_c} + 3m_{J/\psi})$  is used as an input for  $m_{charm}$ , differences from  $m(1S)$  are predictions.
- Our results agree with experiments except for the hyperfine splitting.
- Our hyperfine splitting deviates from BES III experiment(2011) by  $3\sigma$  (4%).  
→ Our error does not include the following systematic errors: scaling violations, dynamical charm quark effects, disconnected loop contributions, QED effects.



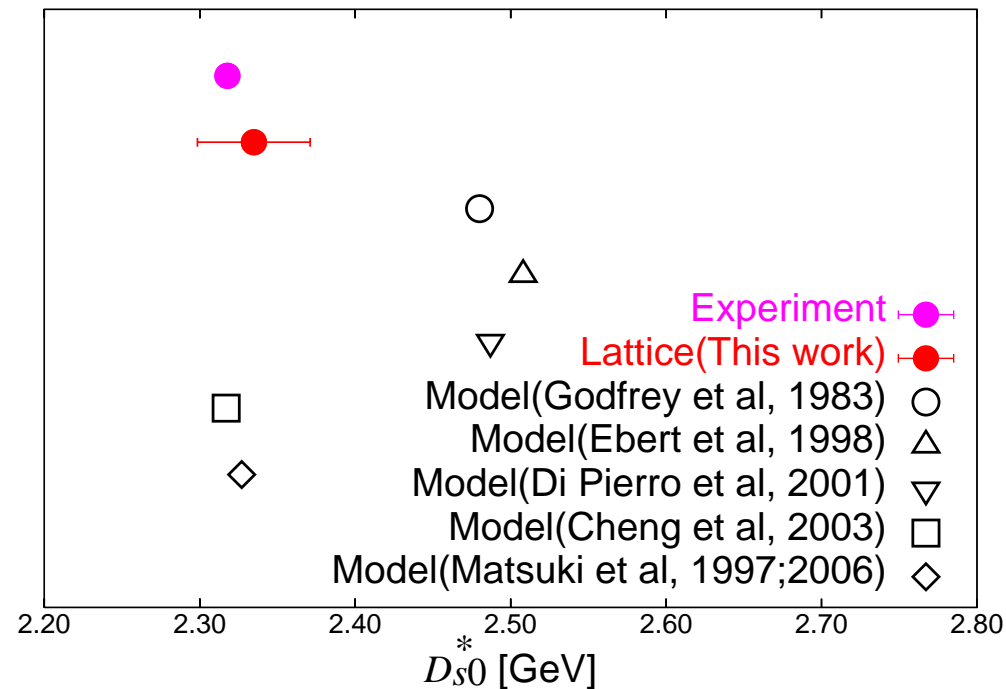
## 3.2 Charm-strange spectrum

- Our calculation reproduces the charm-strange spectrum in  $2\sigma$  level.
- Contaminations to  $m_{D_{s0}^*}$ ,  $m_{D_{s1}}$  from  $DK$  scattering states can be considerably large, which have not been included yet.
- ( $D_{s0}^*$ ,  $D_{s1}$  decays are prohibited in our  $N_f = 2 + 1$  lattice QCD.)



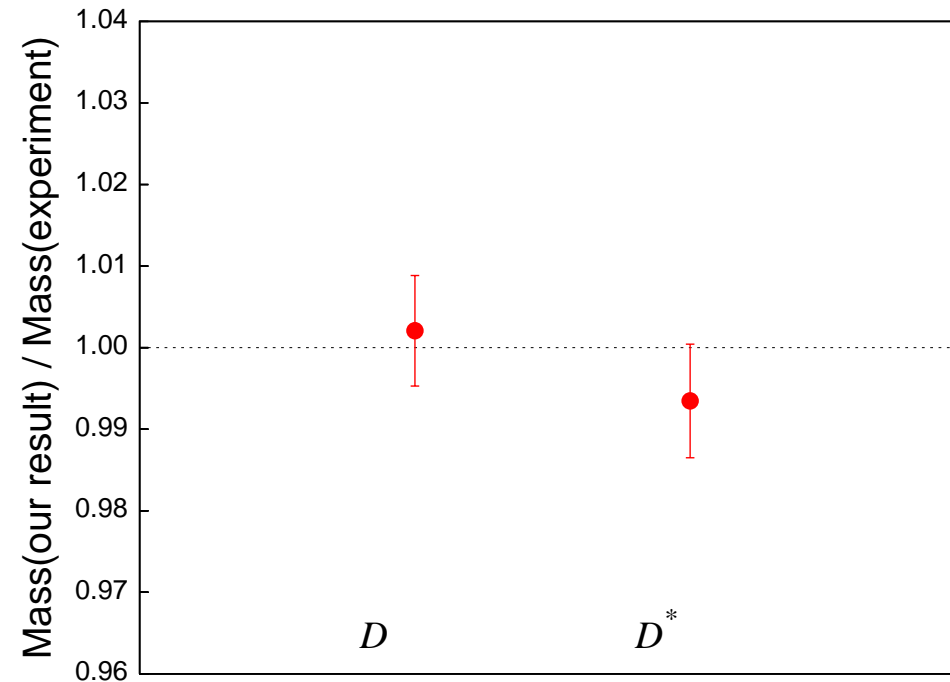
## [Comparison of lattice QCD with non-lattice results]

- Our lattice QCD result using two-quark interpolating operators reproduces the experimental value. No need for multi-quark picture.
- Many models are not good for  $D_{s0}^*$ .
  - ← The standard potential model by Godfrey et al, 1983 fails to reproduce  $D_{s0}^*$  masses.
- A tetra-quark model by Cheng et al, 2003 can fit to experiment.
- A model by Matsuki et al, 1997; 2006 is close to experiment.



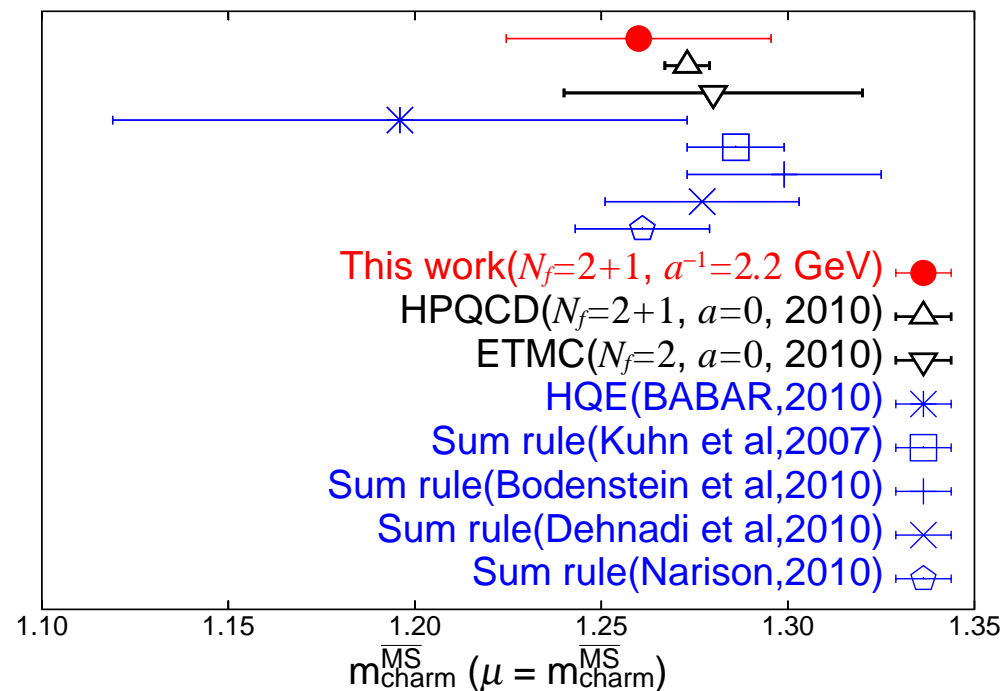
## 3.3 Charm-ud spectrum

- D meson mass spectrums are reproduced well.
- ( $D^*$  decay is prohibited on our lattice of  $L = 3$  fm with  $a^{-1} = 2.2$  GeV.)
- (For unstable particles,  $D_0, D_1$ , more detailed analysis using Lüscher's formula is needed.)



## 3.4 Charm quark mass

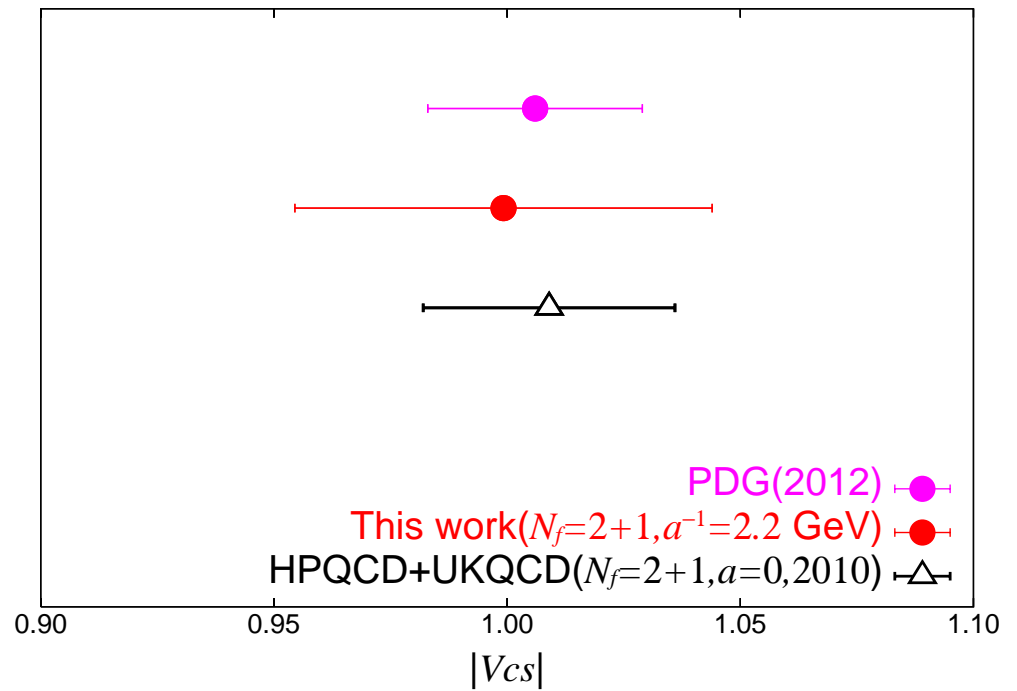
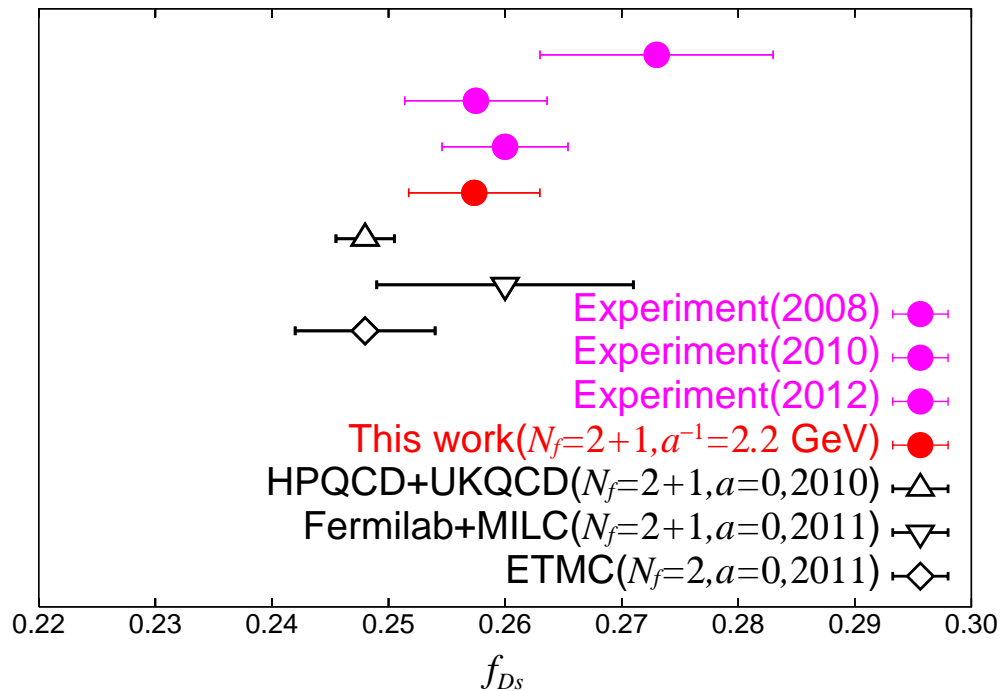
- Charm quark mass is determined from the axial Ward-Takahashi identity.
- Our result is consistent with other lattice and continuum calculations.
- Our systematic error is still large. The main source of our error is the non-perturbative renormalization factors.  
(The renormalization factor is calculated non-perturbatively at the massless point. The mass dependent part is calculated perturbatively.)
- (Charm quark mass is renormalized at  $\mu = 1/a$ , and evolved to  $\mu = m_{charm}^{\overline{MS}}$  using  $N_f = 4$  four-loop beta function.)



## 3.5 Decay constants and CKM matrix elements

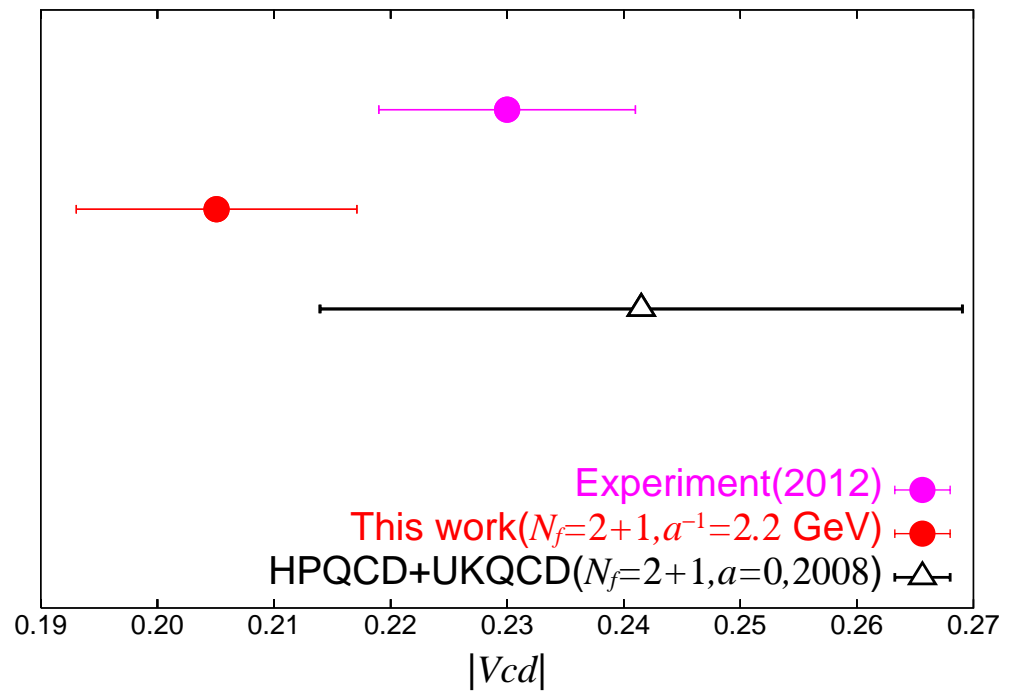
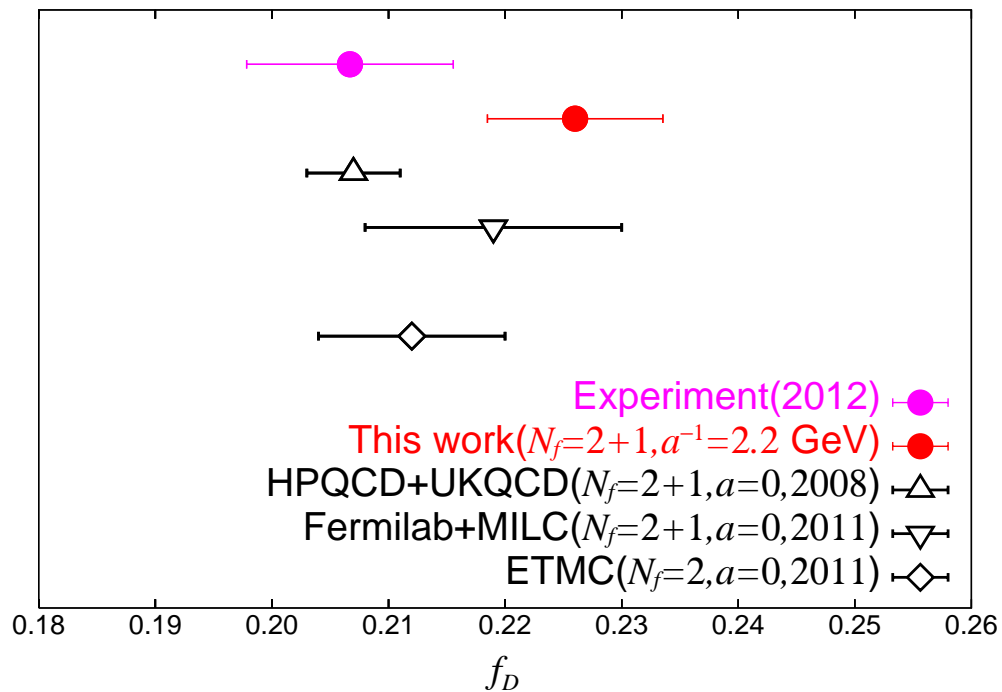
- Our  $f_{D_s}$  agrees with the experimental value and other lattice QCD results.
- CKM matrix elements are extracted from our mass and pseudoscalar decay constant of charmed-strange meson combined with experimental values for the leptonic decay width of charmed-strange mesons.

$$\Gamma(D_s \rightarrow l\nu) = \frac{G_F^2}{8\pi} m_l^2 m_{D_s} f_{D_s}^2 \left(1 - \frac{m_l^2}{m_{D_s}^2}\right)^2 |V_{cs}|^2$$



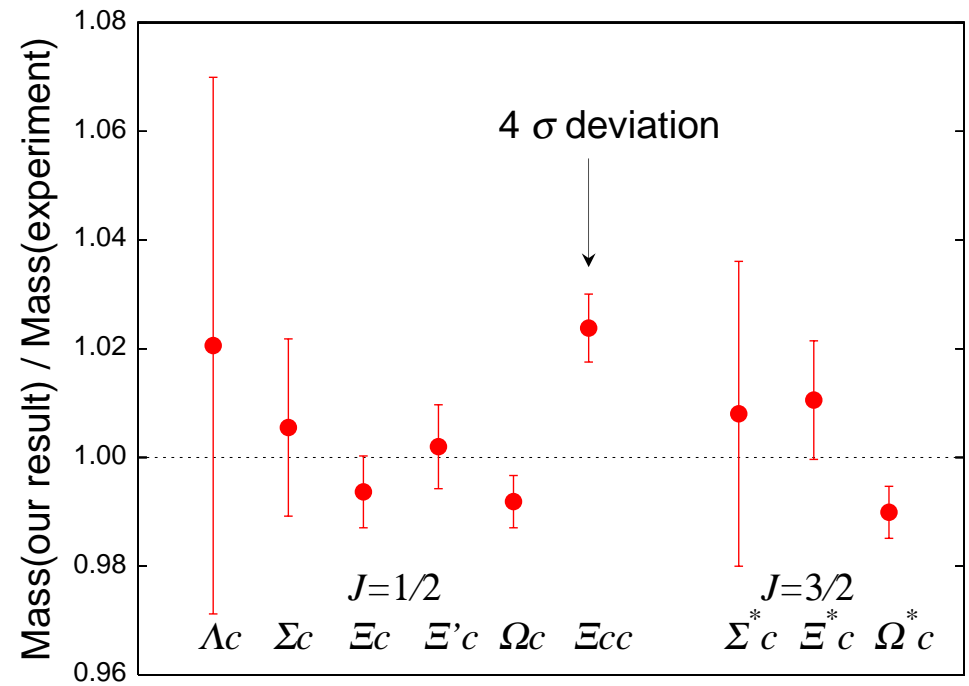
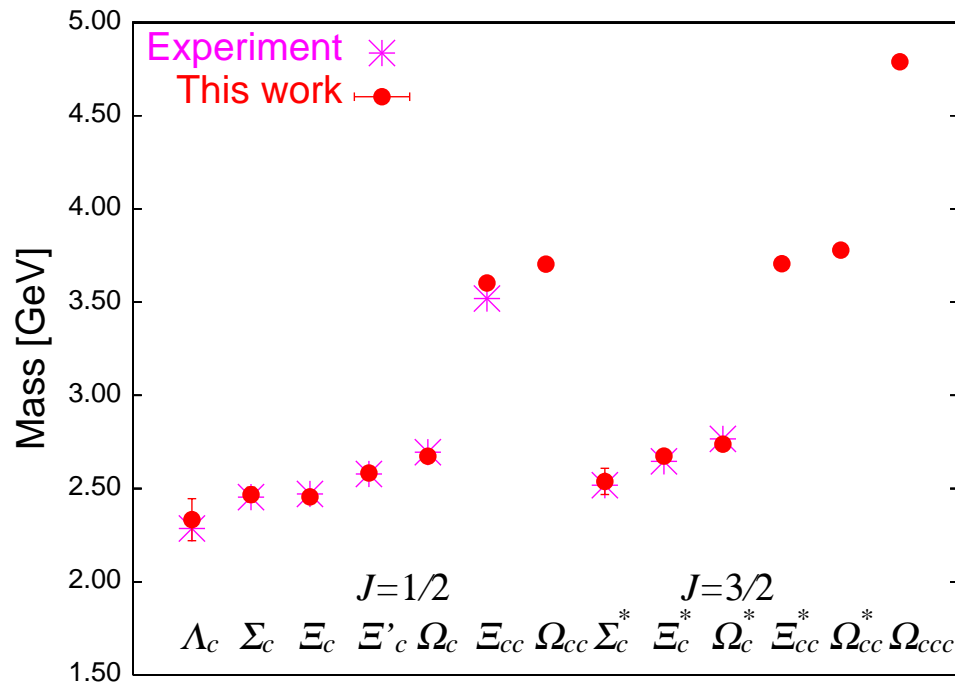
# [Decay constants and CKM matrix elements(continued)]

- Our  $f_D$  is consistent with experiments and other results in  $2\sigma$  level.  
 $\leftarrow$  Our  $f_D$  is  $2.2\sigma$  higher than experiments, and  $2.4\sigma$  higher than HPQCD and UKQCD value.
- Continuum extrapolation of our  $f_D$  is needed for a definite comparison.
- $|V_{cd}|$  is consistent with the other result, because the experimental error of the leptonic decay width is large.



## 3.6 Charmed baryon

- Our results agree with experiments in  $2\sigma$  level, except for  $\Xi_{cc}$ .
  - ◇ Only SELEX(2002,2005) found  $\Xi_{cc} = 3519$  [MeV].
  - ◇ BABAR, BELLE and FOCUS found no evidence for  $\Xi_{cc}$ .  
 →  $\Xi_{cc}$  has been omitted from PDG.

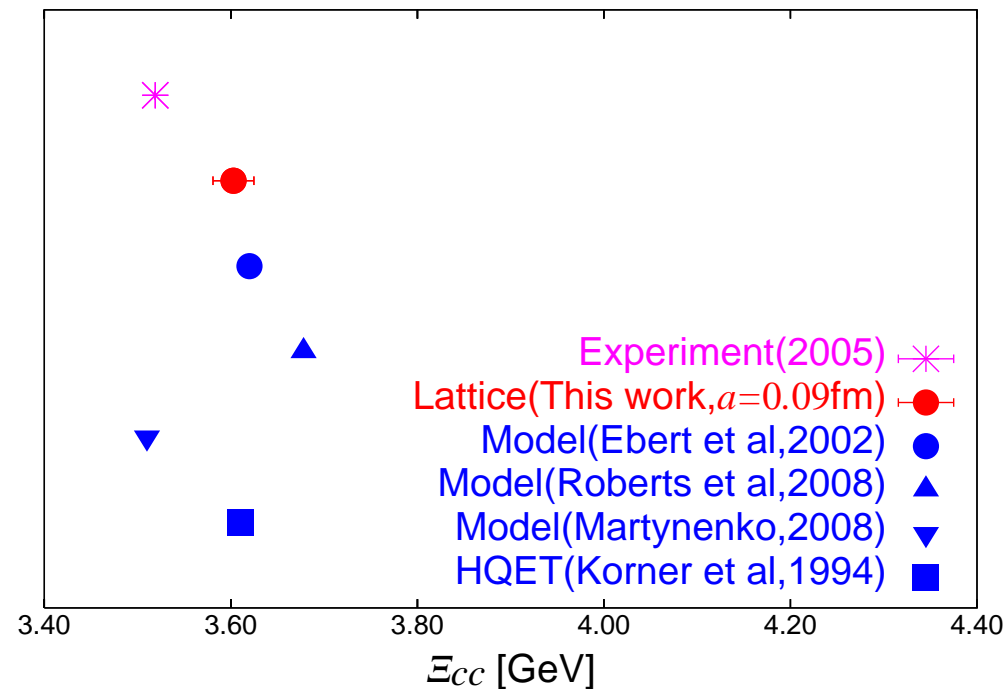




# [Comparison of lattice QCD with non-lattice results]

Typical model calculations are compared with our lattice result.

- For  $\Xi_{cc}$ , many models are close to our result.
- Many models and our result give  $\Xi_{cc}$  mass higher than SELEX experiment by around 100 MeV.



# 4 Summary

We performed  $N_f = 2 + 1$  full QCD simulation of the charm quark system on the physical point at  $a^{-1} = 2.2$  GeV.

- Our calculation reproduces meson mass spectrums of the ground states except for hyperfine splittings.
  - ◇ Our data of the charmonium hyperfine splitting is  $3\sigma$  smaller than experiments.
    - ← Possible origins of the discrepancy are  $O(a)$  effects in our relativistic heavy quark action, dynamical charm quark effects, disconnected loop contributions, QED effects.
- Our results for charm quark mass and CKM matrix elements are presented.
- Our calculation reproduces charmed baryon spectrum except for  $\Xi_{cc}$ .
  - ◇ Our data of  $\Xi_{cc}$  shows a significant deviation from the experimental value of SELEX group.

[Future works]

- We are going to a finer lattice ( $a^{-1} = 3$  GeV) to take a continuum limit.
- Excited states of charmonium such as  $X, Y, Z$ , separating  $D\bar{D}$  contamination.

[New computers]

Machine	Speed [PFlops]
K-computer@RIKEN,AICS	11
BlueGene/Q@KEK	1.3
HA-PACS@Univ. of Tsukuba	0.8
PACS-CS@Univ. of Tsukuba	0.01
CP-PACS@Univ. of Tsukuba	0.001



# Appendix